

Strongly Coupled Theories of the TeV Scale

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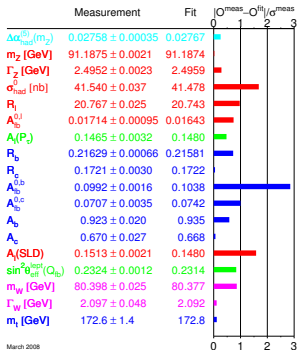
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Outline

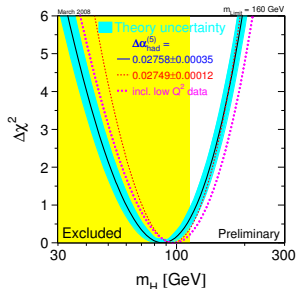
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The great success of the SM

The Standard Model success



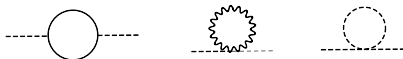
The SM Higgs is light



$m_h < 190 \text{ GeV} @ 95\% \text{ C.L.}$

The Stability of the Weak Scale

But if Higgs elementary scalar quantum corrections drive m_h up



$$\delta m_h^2 \sim \frac{c^2}{16\pi^2} \Lambda^2$$

- We need $\Rightarrow m_h \lesssim 1 \text{ TeV}$
- But if $\Lambda \rightarrow M_P \sim 10^{19} \text{ GeV}$, unnatural

\Rightarrow Gauge Hierarchy Problem

Mechanism to Stabilize the Weak Scale

Need a mechanism to keep v and M_P separate.

New physics at $\Lambda \sim 1$ TeV is:

Weakly Coupled

- SM with a light Higgs
- SUSY (MSSM, NMSSM, Folded, ...)
- Little Higgs, Twin Higgs
- LED, UED

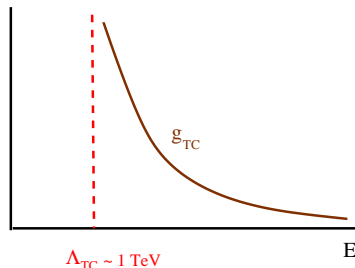
Strongly Coupled

- Technicolor, Walking Technicolor
- Topcolor, Top See Saw
- Composite Higgs
- Randall-Sundrum

Strong Dynamics at the TeV Scale

Analogy with QCD:

- New Strong Interaction: Technicolor
- Strong at $M_W \ll M_P$
- Breaks Electroweak symmetry: $\langle \bar{F}F \rangle \neq 0$



Strong Dynamics at the TeV Scale

TC Problem I: Fermion Masses

- Need Extended Technicolor (ETC):

$$G_{\text{ETC}} \supset G_{\text{TC}} \text{ with } \Lambda_{\text{ETC}} \gg \Lambda_{\text{TC}}$$

$$m_f \simeq \frac{g_{\text{ETC}}^2}{\Lambda_{\text{ETC}}^2} \langle \bar{F}F \rangle$$

- Needs Walking (Walking TC):
For heavier fermions (e.g. m_c, m_τ) needs to enhance $\langle \bar{F}F \rangle$
- But m_t requires Λ_{ETC} too low
 \Rightarrow Topcolor, Top See Saw (Hill, Dobrescu)

Strong Dynamics at the TeV Scale

TC Problems II: Electroweak Precision Constraints



- E.g. Techni-fermions give

$$S \sim \frac{N}{6\pi}$$

- But

$$S^{\text{exp.}} \leq 0.1$$

Strong Dynamics and AdS/CFT

AdS/CFT Correspondence (Maldacena):

- Originally:

$\text{AdS}_5 \times S^5$ String Theory $\leftrightarrow \mathcal{N} = 4$ 4D SU(N) Theory (CFT)

- In general: Assume 5D theory in $\text{AdS}_5 \leftrightarrow$ 4D CFT
(Arkani-Hamed, Porrati, Randall)

- Need

$$g^2 N \gg 1$$

to ignore string corrections.

- \Rightarrow Holographic dual is 4D strongly coupled theory

Strong Dynamics from a Slice of AdS₅

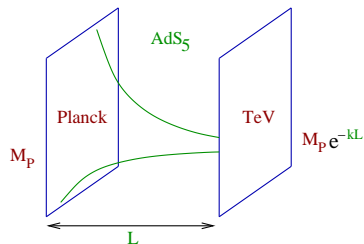
Ingredients to build Strongly Coupled Theories in AdS₅

- UV cutoff in the 4D Theory \leftrightarrow UV (“Planck”) boundary
- Break 4D Conformal Invariance in the IR \leftrightarrow IR boundary
- 4D Strongly Coupled Gauge Theory described by 5D Weakly Coupled Theory

Solving the Hierarchy Problem in AdS₅

Metric in extra dimension \Rightarrow small energy scale from M_P
(Randall-Sundrum)

$$ds^2 = e^{-2\kappa|y|} \eta^{\mu\nu} dx_\mu dx_\nu - dy^2$$



Corrections to m_h OK
If Higgs close to TeV brane

Need Higgs IR localization

Bulk AdS₅ Models

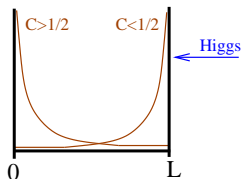
Allowing Gauge and Matter fields in 5D bulk
(Gherghetta-Pomarol, Grossman-Neubert)

- Avoid effects of Higher Dimensional Operators only suppressed by IR/TeV scale
- Natural Models of Flavor:
Zero-mode fermion localization ↔ fermion mass

$$M_f^{(5D)} = ck, \quad c \simeq O(1)$$

Fermion Masses in Bulk RS Models

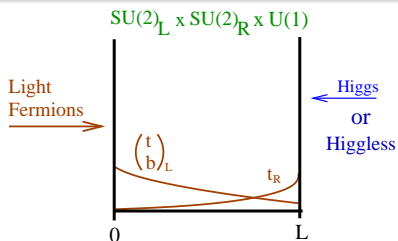
- $O(1)$ flavor breaking in bulk can generate fermion mass hierarchy:



TeV localization \rightarrow larger Yukawas,
Planck localization \rightarrow suppressed Yukawas.

- Heavier fermions couple stronger to gauge KK modes:
 - $G^{(1)} \rightarrow t\bar{t}$ dominates
 - Tree-level flavor violation

The Bulk RS Picture



Models of
 EWSB *and* Flavor

- EWPC: T OK, but $S \simeq N/\pi$ at tree-level

$$M_{KK} \gtrsim (2 - 3) \text{ TeV}$$

- $Z \rightarrow \bar{b}b$ require discrete symmetry ($L \leftrightarrow R$)
 (Agashe, Contino, Da Rold, Pomarol)
- Potentially important bounds and/or effects
 from flavor violation

Dynamical Origin of the Higgs Sector

What localizes the Higgs to/near the IR/TeV brane ?

- Gauge-Higgs Unification
- Zero-mode Fermion Condensation
- Higgsless

Gauge-Higgs Unification in AdS₅

Extract H from a 5D gauge field

- $A_M^a \rightarrow (A_\mu^a, A_5^a)$
- Bulk Gauge group:
 - $G_{\text{bulk}} \supset G_{\text{SM}}$
 - Needs to include custodial symmetry.
- Minimal model (Agashe, Contino, Pomarol):
 $SO(5) \times U(1) \rightarrow SO(4) \times U(1)$ by BCs
- Higgs is **4** of $SO(4)$: 4 d.o.f. \leftrightarrow complex $SU(2)_L$ doublet
- H localized near TeV brane.
 m_h protected by bulk gauge symmetry (H pseudo-NGB)

Fermion Condensation

Fourth-Generation Condensation in AdS₅ (G.B. Da Rold)

- Fourth Generation in the AdS₅ bulk
- Choose zero-mode fermions IR localized \Rightarrow strongly coupled to KK gauge bosons



E.g. KK gluon exchange $\longrightarrow \langle \bar{U}U \rangle \neq 0$

- EWSB
- $m_U^{(0)} \sim (600 - 700) \text{ GeV}$ (ala Bardeen-Hill-Lindner)
- Heavy Higgs: $m_h \simeq (600 - 900) \text{ GeV}$

Fermion Condensation (cont.)

- All other fermion masses: Bulk higher dimensional operators

$$\frac{C^{ijkl}}{M_P^3} \bar{\Psi}_L^i(x, y) \Psi_R^j(x, y) \bar{\Psi}_R^k(x, y) \Psi_L^\ell(x, y)$$

- Phenomenology dominated by 4th generation
 - $V^{(1)} \rightarrow \bar{U}U$ (broader KK gauge bosons)
 - Flavor physics: E.g. new sources of CPV in mixing, ...
 - Additional contributions to S, T

Higgsless Models

Higgsless RS Bulk Models (Csaki, Grojean, Murayama, Pilo, Terning)

- Boundary Condition breaking

$$SU(2)_L \times SU(2)_R \times U(1)_X \rightarrow U(1)_{EM}$$

- IR localized mass terms \Rightarrow fermion masses
- Kaluza-Klein modes of gauge fields unitarize amplitudes.
 \Rightarrow KK modes “light”: $M_{KK} \lesssim 1 \text{ TeV}$
- Phenomenology in the Gauge boson sector:
 - $V_L V_L$ scattering
 - Sum Rules
- Corresponds to Walking Technicolor Models

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- If LHC reveals Strongly Coupled TeV scale
⇒ Model Building in AdS₅ should be a useful tool