

At the Edge of Discovery

Beyond the Standard Model of Particle Physics

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Outline

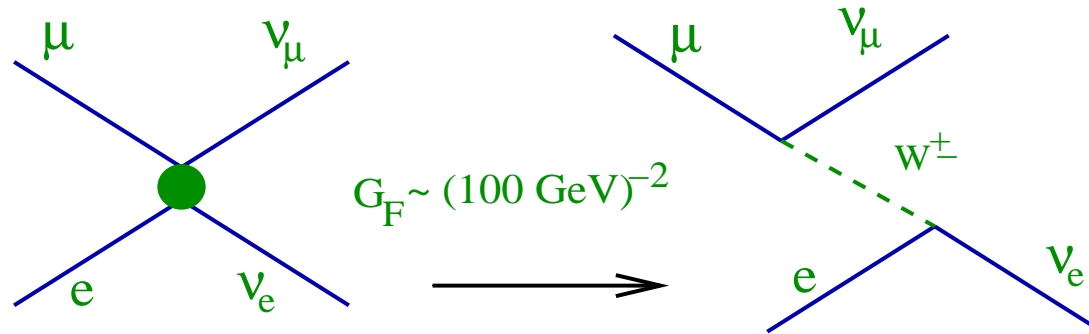
- * The Standard Model: a successful description.
- * The hierarchy problem. Why do we believe that the TeV scale will reveal new dynamics ?
- * Physics beyond the Standard Model: understanding M_W/M_{Planck} .
Old and New approaches to the TeV scale.
 - Supersymmetry
 - Composite scalar sectors (Technicolor, Topcolor, Little Higgs, ...)
 - Extra Dimensions (Large, Universal, Warped)
- * Conclusions/Outlook

The Road to the Standard Model

- Quantum Electrodynamics (QED):
 - Successful description of Electromagnetism as a quantum field theory.
 - Local Gauge Symmetry: Theory invariant under local $U(1)$ transformations.
 - $\Rightarrow \gamma$'s are the force carriers, and they are **massless**.
 - Tested to great precision!! (Lamb shift, $(g - 2)$, \dots).
- Weak Interactions:
 - Responsible for nuclear β decay: $n \longrightarrow p^+ + e^- + \bar{\nu}_e$.
 - Short range! \Rightarrow carriers are not massless!
 - Inconsistent with Gauge Theories (e.g. QED)?

The Road to the Standard Model

- Weak Interactions: Mediated by massive vector bosons



- Gauge Symmetry now is $SU(2)_L \times U(1)_Y$:
Discovery of neutral currents in 1970's \Rightarrow massive (W^\pm, Z^0), plus a massless γ .
- Massive gauge boson \Rightarrow broken gauge symmetry !

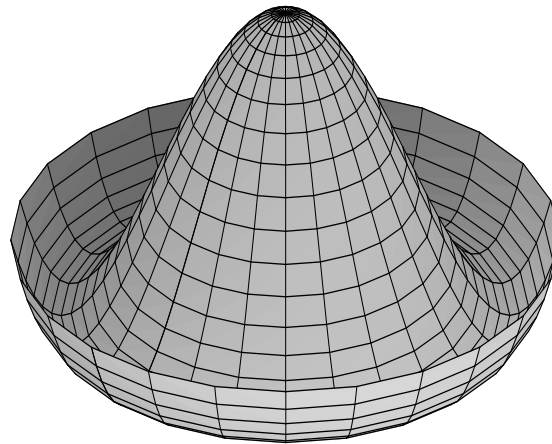
$$\boxed{m^2 A_\mu A^\mu} \text{ not invariant under } \boxed{A_\mu(x) \rightarrow A_\mu(x) - \frac{1}{g} \partial_\mu \alpha(x)}$$

Theory is inconsistent at high energies ? (Non-renormalizable)

Spontaneous Symmetry Breaking

Symmetry is not broken but **hidden**.

True vacuum not symmetric \Rightarrow spontaneously broken symmetry.



\Rightarrow Massless excitations in the “symmetric” direction

\longrightarrow Goldstone bosons

Superconductivity

Superconductor: Electromagnetic gauge invariance is spontaneously broken. Material obeys

$$\mathcal{E}_{\text{normal}} = \mathcal{E}_{\text{SC}} + \Delta$$

Superconducting vacuum breaks $U(1)_{\text{EM}} \Rightarrow$ Goldstone boson ϕ such that

$$A_\mu = \partial_\mu \phi \text{ minimizes energy.}$$

\Rightarrow Magnetic Field vanishes inside !! \rightarrow flux exclusion (Meissner effect.)

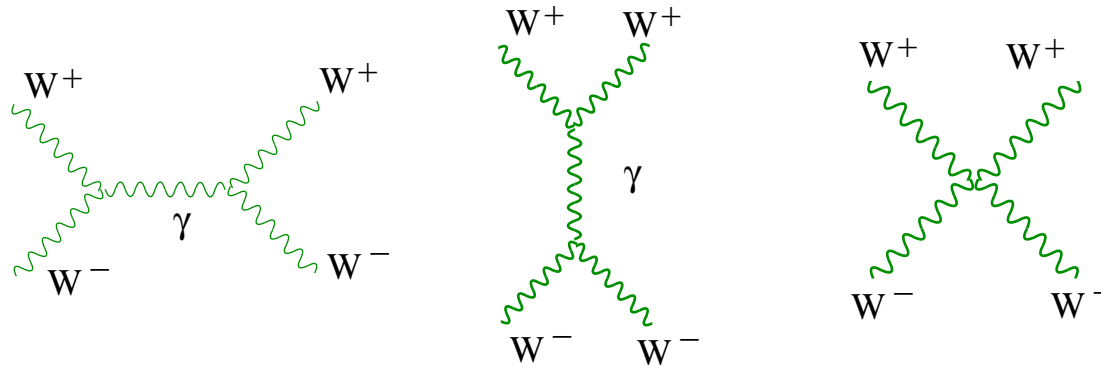
\Rightarrow EM interactions become short range \Rightarrow “Massive” Photon

The Higgs Mechanism

- Need to spontaneously break $SU(2)_L \times U(1)_Y$ to $U(1)_{EM}$. Only **photon** remains massless.
- Higgs: Goldstone bosons are “eaten” by (W^\pm, Z^0) .
 $\Rightarrow (m_W, m_Z) \Rightarrow$ **Weak Interactions** become short range.
- In the Standard Model, one scalar particle (**The Higgs Boson**) remains in the spectrum.
- In the case of Superconductivity, **Cooper pairs** form a condensate that breaks $U(1)_{EM}$.
- Assume the universe is a “superconductor”. Condensate breaks the electroweak symmetry.

The Higgs Mechanism

- At what energy scale should the condensate occur ?
- E.g. in W^+W^- scattering



Amplitudes grow like $\sim s/(M_W^2)$

They would violate unitarity at $\Lambda \sim 1 \text{ TeV}$

The Standard Model

- Introduce $\Phi = \begin{pmatrix} \phi^+ \\ \phi^0 \end{pmatrix}$, the Higgs doublet:

$$\mathcal{L} = (D_\mu \Phi)^\dagger (D^\mu \Phi) - V(\Phi) \quad SU(2) \times U(1) \text{ invariant}$$

- Interactions with gauge bosons in $D_\mu = \partial_\mu - g\vec{W}_\mu - g'B_\mu$

- $V(\Phi) \Rightarrow$ vacuum not invariant: $\langle \Phi \rangle = \begin{pmatrix} 0 \\ v \end{pmatrix}$

\Rightarrow masses for 3 out of the 4 gauge bosons.

- $(\vec{W}, B) \longrightarrow (W^\pm, Z^0)$, plus a massless γ

$$M_W, M_Z \neq 0, \quad m_\gamma = 0$$

The Standard Model

- To avoid unitarity violation in cross sections

$$\Rightarrow m_h < 1 \text{ TeV}$$

The Success of the Standard Model

- Precision Tests (1990's):

- Millions of Z^0 's produced at LEP I via

$$e^+e^- \longrightarrow Z^0 \longrightarrow f\bar{f}$$

- Non-abelian couplings tested at LEP II in

$$e^+e^- \longrightarrow W^+W^-$$

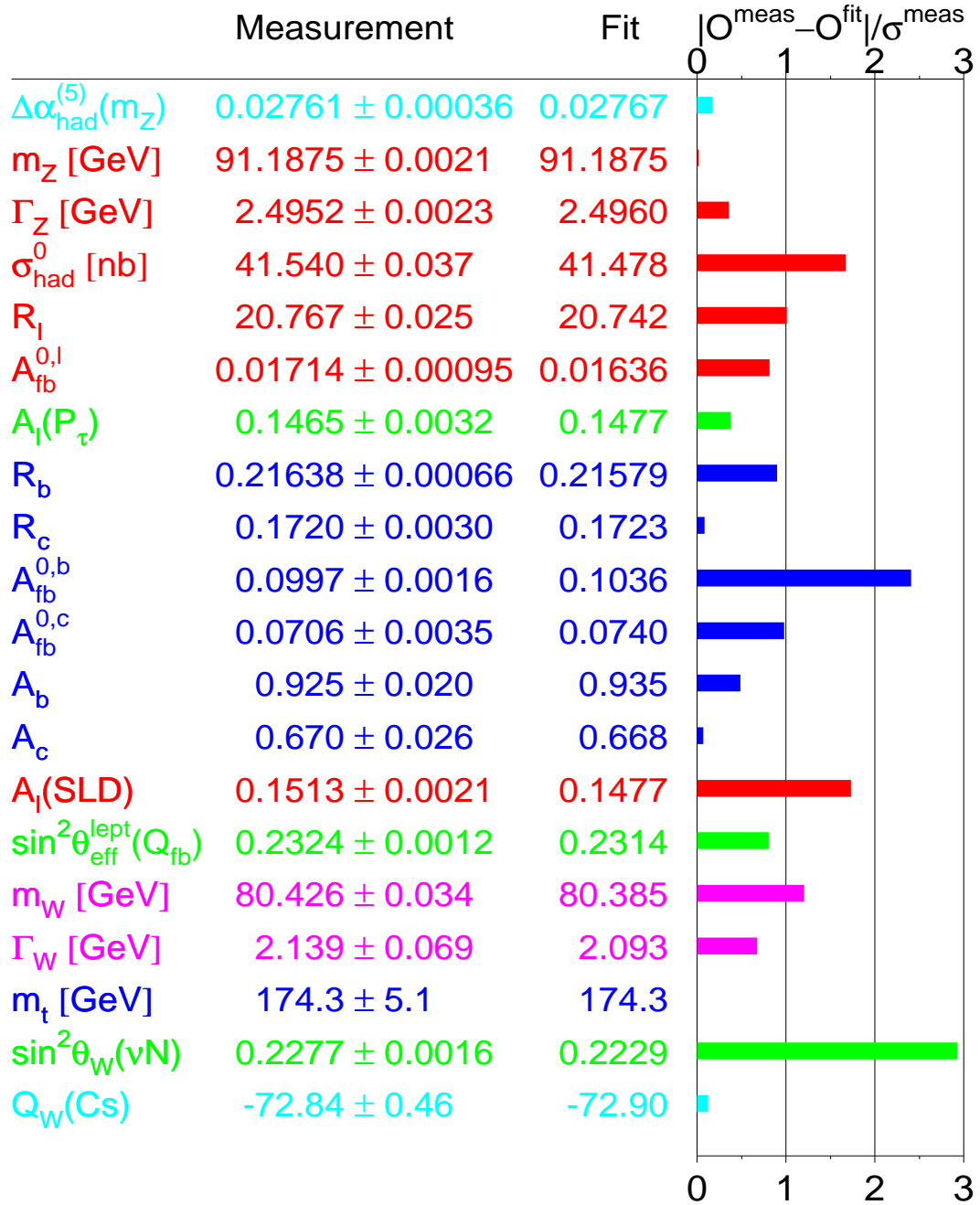
- Also precision νN and Atomic Parity Violation experiments.

- Input most precise measurements

$$M_Z, \alpha \text{ and } G_F \text{ (where } G_F \text{ is mostly from } \mu^\pm \rightarrow \nu_\mu e^\pm \nu_e)$$

\implies Standard Model fit to all data

Summer 2003



The Limitations of the Standard Model

- The Hierarchy Problem:

Why is the **weak scale** ($\lesssim 10^3$ GeV), so much smaller than the gravitational scale ($M_{\text{P}} \simeq 10^{19}$ GeV) ?

- In the **Standard Model** the weak scale is **not stable** !!!

$$\Delta m_h^2 = \text{---} \overset{m_t}{\bullet} \text{---} \text{---} \text{---} \overset{m_t}{\bullet} \text{---} \text{---} + \dots$$

$$\Rightarrow \Delta m_h^2 \simeq 3\Lambda^2 (2M_W^2 + M_Z^2 + m_h^2 - 4m_t^2) / (32\pi^2 v^2)$$

But we need $\Lambda < 1$ TeV

Another Ultraviolet Catastrophe ?

- Need new physics at the **TeV** scale to stabilize the weak scale.
 - Additional states cancel divergences due to symmetries (e.g. **Supersymmetry**)
 - Higgs is composite and “comes apart” at scale Λ .
 - $(\text{TeV})^{-1}$ is the size of Extra Spatial dimensions

⋮

- The analogy with Superconductivity: The **Higgs sector of the Standard Model** is the “Landau-Ginzburg” theory. The BCS dynamics must appear at the **TeV** scale

The Limitations of the Standard Model

- What is the origin of Fermion masses ?:

In the Standard Model, *ad hoc* couplings of Higgs to fermions are adjusted to obtain

$$(m_e)/(m_t) \sim 10^{-6}, \quad m_\nu \lesssim 1 \text{ eV}$$

- Do interactions Unify at high energies ?

$$SU(3) \times SU(2)_L \times U(1) \longrightarrow G ?$$

- What is the **Dark Matter** ?
- What is the **Dark Energy** ?

⋮

Supersymmetry at the Weak Scale

- In the Supersymmetric limit quadratic div. from superpartners cancel the ones from the SM particles:

$$\Delta m_h^2 = \text{---} \text{---} \text{---} + \text{---} \text{---} \text{---} + \dots$$

The diagram shows two Feynman diagrams representing the top quark loop contribution to the Higgs mass. The first diagram is a solid blue circle with 't' at the top and bottom, and 'm_t' at the left and right vertices. The second diagram is a dashed blue circle with 't-tilde' at the top and 'm_t^2' at the bottom vertex. Green dashed lines represent external Higgs lines.

- Remaining div. is logarithmic

$$\Rightarrow \Delta m_h^2 \simeq m_{\text{soft}}^2 c / (16\pi^2) \ln(\Lambda_{\text{SUSY}} / m_{\text{soft}}) + \dots$$

- Weak scale SUSY also helps SU(5) coupling unification, results in radiative EWSB because m_t is large, has DM candidates, \dots

Supersymmetry at the Weak Scale - The Problems

- In **SUSY**, at tree-level $m_h < m_Z$.
- But we know $m_h^{\text{exp.}} \geq 114$ GeV. The logarithmic radiative corrections must lift m_h .
- $\Rightarrow m_{\text{soft}}$ (e.g. squark masses) must be high (\simeq TeV)
- This results in (*severe* ?) **fine tuning** of radiative EWSB !
- Extensions of the MSSM: NMSSM, Extra Dimensions, Split SUSY (still unnatural), \dots

A Composite Higgs Sector

- Vacuum is “Electroweak Superconductor”. New gauge interactions at the TeV scale \implies Fermion condensate (Cooper pairs):

$\langle \bar{F}F \rangle \neq 0$ breaks electroweak symmetry $\rightarrow U(1)_{EM}$

- In *QCD* condensate breaks chiral symmetry
 \implies “constituent” quark masses $m_q \sim \Lambda_{QCD} \sim 300$ MeV.
- “Technicolor” becomes strong at $\Lambda \sim 1$ TeV
 \implies 3 Goldstone bosons “eaten” by W^\pm and Z^0
 $\implies M_W, M_Z$.
- Higgs dissolves into its constituents at $\Lambda \sim 1$ TeV
 \implies No hierarchy problem .

A Composite Higgs Sector

- Topcolor: $m_t \sim 175 \text{ GeV} \sim v$...suggestive ... (C. Hill)

$$m_t \sim O(\Lambda) \rightarrow \text{is a constituent mass ?}$$

- New gauge interaction couples strongly to the top quark
 $\Rightarrow \langle \bar{t}_L t_R \rangle \neq 0 \Rightarrow m_t$

But $\begin{pmatrix} t \\ b \end{pmatrix}_L \Rightarrow SU(2)_L \times U(1)_Y$ broken to EM.

- Higgs sector: Higgs boson (mostly) a $\bar{t}t$ composite.
 \Rightarrow couples strongly to the third generation quarks t, b .
 \Rightarrow Flavor violation at tree level (FCNCs).

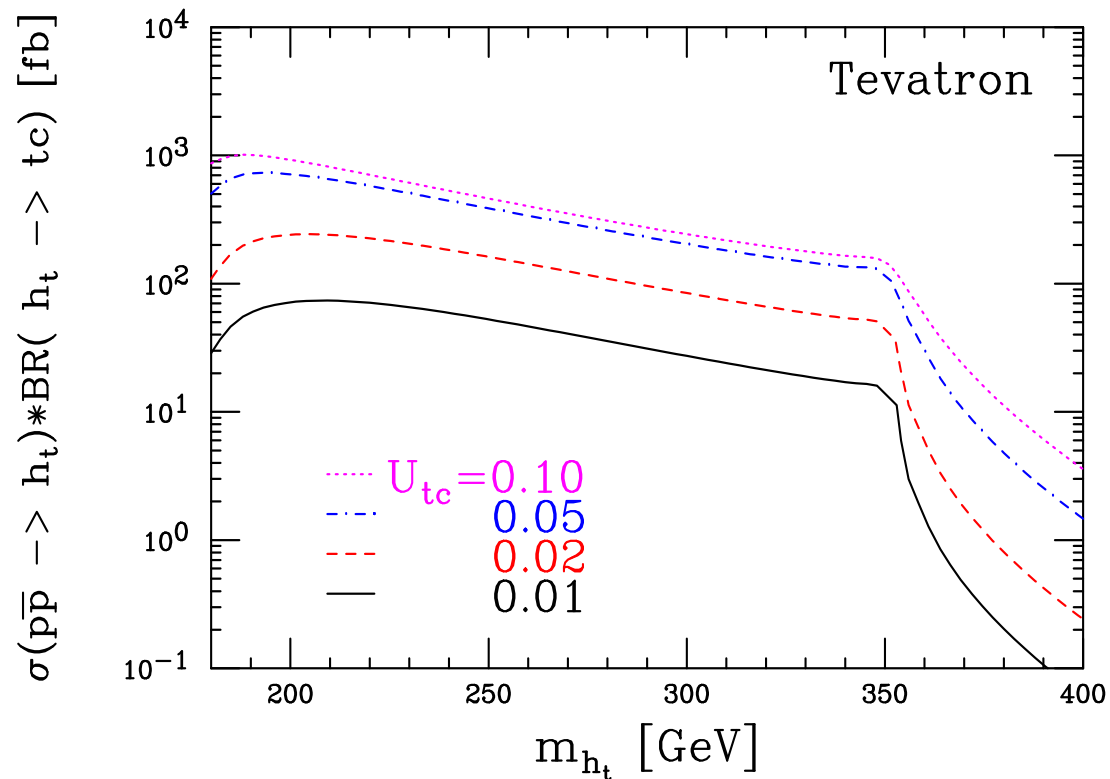
Tight constraints to model building from precision measurements and weak decays of heavy quarks (G. Buchalla, G.B., D. Kominis, C. Hill).

Experimental Signals of a Composite Higgs Sector

The “top-Higgs” h_t decays mostly to flavor violating channel $h_t \rightarrow tq$.

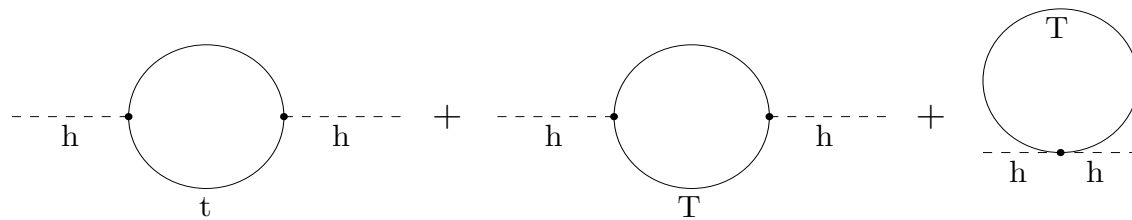
\Rightarrow It can be observed at Tevatron RunII and LHC !! (G.B.).

$$p\bar{p} \rightarrow \boxed{gg \longrightarrow h_t \longrightarrow tq} \rightarrow bl\nu_e + \text{jet}$$



A Little Help from a Little Higgs

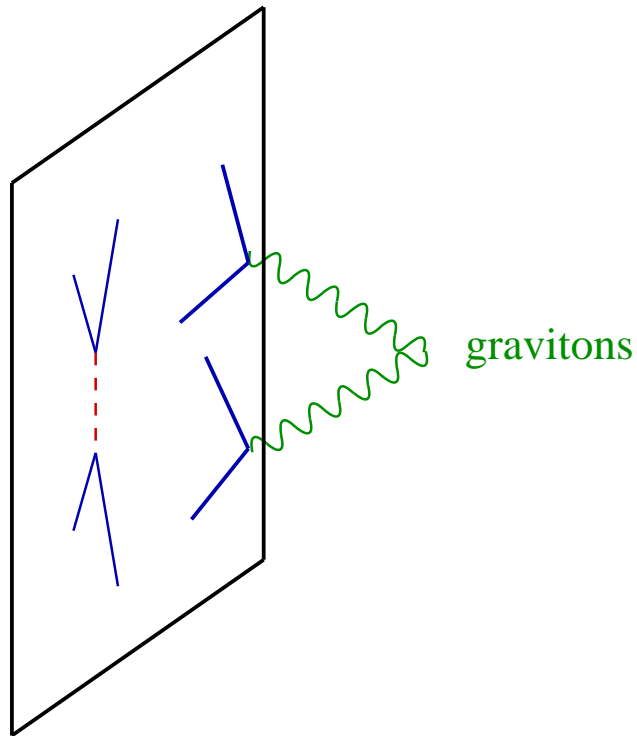
- Theories of composite Higgs highly constrained by precision data. Strong dynamics at $\Lambda \sim \text{TeV}$ scale is difficult. Raising Λ reintroduces “little hierarchy”.
- Can raise Λ as long as we cancel Δm_h^2 at the TeV scale. Introduce new global (and gauged) symmetries \Rightarrow cancellations from loops (Arkani-Hamed, Cohen, Georgi).



- New particles at the TeV scale: 4 new gauge bosons, 1 or more new heavy quarks.

Extra Dimensions and the Hierarchy Problem

- Assume space has $3 + n$ dimensions.
- The extra n dimensions are compact and with radius R .
- All particles are confined to a **3-dimensional** slice (“brane”).
- Gravity propagates in *all* $3 + n$ dimensions.



Large Extra Dimensions

(Arkani-Hamed, Dimopoulos, Dvali '98)

- Gravity appears weak ($M_P \ll M_W$), because it propagates in large extra dimensions... Its strength is diluted by the volume of the n extra dimensions:

$$V(r) \sim \begin{cases} \frac{m_1 m_2}{M_*^{n+2}} \frac{1}{r^n} & (r \ll R) , \\ \frac{m_1 m_2}{M_*^{n+2}} \frac{1}{R^n} & (r \gg R) , \end{cases}$$

- Fundamental scale is $M_* \sim M_W$, not M_P

$$M_P^2 \sim M_*^{n+2} R^n$$

- There is no hierarchy problem!

Large Extra Dimensions

- If we require $M_* = 1$ TeV:

$$R \sim 2 \cdot 10^{-17} 10^{\frac{32}{n}} \text{ cm}$$

- $n = 1 \implies R = 10^8$ Km. Already excluded!
- $n = 2 \implies R \simeq 2$ mm. Barely allowed by current gravity experiments.
- $n > 2 \implies R < 10^{-6}$ mm. This is fine.

- Compact extra dimensions \implies graviton excitations (Kaluza-Klein)
Mass gap $\Delta m \sim 1/R$ E.g. for

$$n = 2 \longrightarrow \Delta m = 10^{-3} \text{ eV.}$$

$$n = 3 \longrightarrow \Delta m = 100 \text{ eV.}$$

⋮

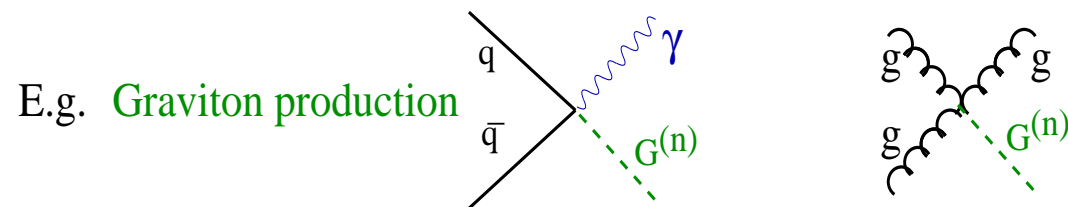
$$n = 7 \longrightarrow \Delta m = 100 \text{ MeV.}$$

Large Extra Dimensions

- Individual KK graviton couplings gravitationally suppressed ($\sim 1/M_P$).
- But for $E \gg 1/R \rightarrow$ sum of KK mode results in

$$\sigma \sim \frac{E^n}{M_*^{n+2}}.$$

- Collider Processes:



Individual graviton decay rates $\sim 1/M_P^2$, $\Rightarrow \cancel{E}_T$ signals at colliders.
Bounds on M_* from LEP and Tevatron (see next).

Large Extra Dimensions

- Virtual KK Graviton Exchange: Induces operators

$$\text{dim} - 6 \sim [\bar{f}\gamma_\mu\gamma_5 f]^2$$

$$\text{dim} - 8 \sim T_{\mu\nu}T^{\mu\nu}$$

- **dim-8** operator leads to $f\bar{f} \rightarrow \gamma\gamma, ZZ, \dots$, pure d-wave.
- Bounds on M_* from these ops. are in the (1-10's) TeV range.

- Black Holes at the LHC? (Giddings, Thomas '02)

- Astrophysical Constraints

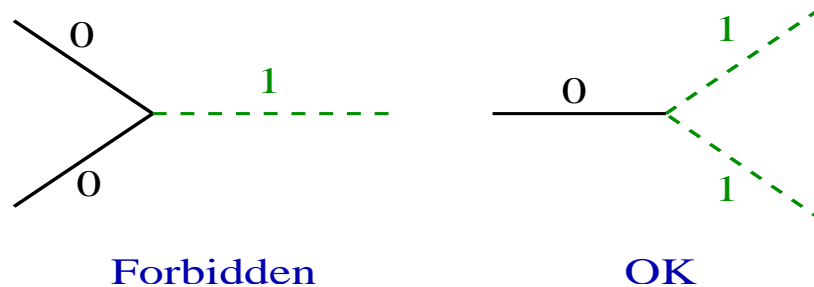
Supernova Cooling: Graviton emission carries too much energy, cooling SN too fast. For $n = 2$, $M_* > (10 - 100)$ TeV.

(Cullen, Perelstein '99, Hall, Smith '99).

Universal Extra Dimensions

(Appelquist, Cheng, Dobrescu '01)

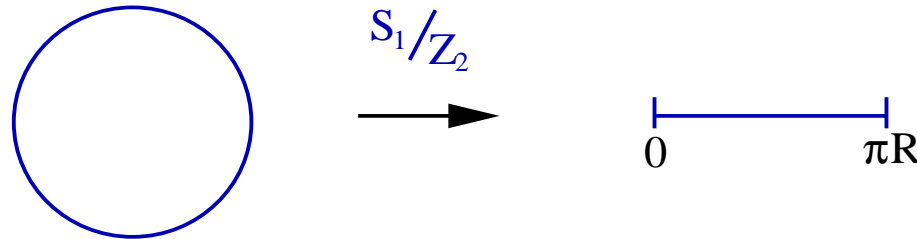
- If some SM fields propagate in the bulk $\Rightarrow 1/R \gtrsim 1 \text{ TeV}$.
- But if we assume *all* fields can propagate in the extra dimensions. What is the allowed R ?
- Momentum conservation in the extra dimensions \Rightarrow KK-number conservation E.g.



\Rightarrow KK excitations must be pair produced, direct bounds on $1/R$ are lower.

Universal Extra Dimensions

- Orbifold Compactification: Allows to have chiral fermions



- It breaks $\text{KK-number} \longrightarrow \text{KK-parity}$

\Rightarrow Lightest KK Particle (LKP) is **stable** \rightarrow **Dark Matter candidate**

- Direct and EW constraints:

$$1/R \gtrsim 300 \text{ GeV for 5D}$$

$$1/R \gtrsim (400 - 600) \text{ GeV for 6D}$$

Universal Extra Dimensions

- Light KK modes \Rightarrow large cross sections.
- But, almost degenerate KK levels \Rightarrow little energy release:
E.g. $q\bar{q} \rightarrow Q_1 Q_1 \rightarrow Z_1 Z_1 + \cancel{E_T} \rightarrow 4\ell + \cancel{E_T}$ (Cheng, Matchev, Schmaltz '02).
- 6D case well motivated:
 - Proton Stability (Appelquist, Dobrescu, Pontón, Yee '01).
 - Three generations (Dobrescu, Poppitz '01)

Signals somewhat different from 5D (Burdman, Dobrescu, Pontón '04)

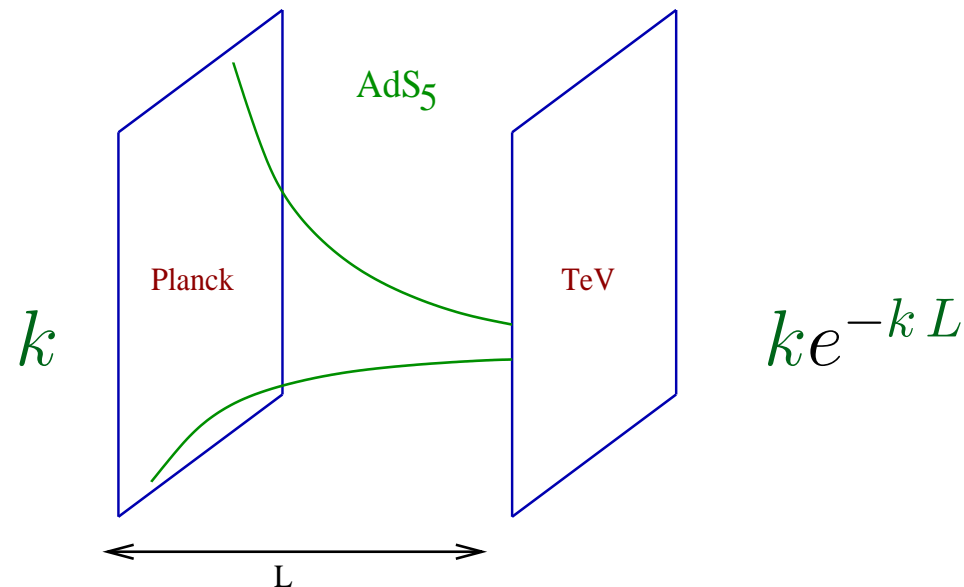
E.g. Second KK level *cannot* decay to two level-1 KK modes:

$$m_2 \simeq \sqrt{2}/R.$$

Can decay to two zero-modes through KK-number violating brane kinetic terms.

Warped Extra Dimensions

- One compact extra dimension. Non-trivial metric induces small energy scale from Planck scale! (L. Randall, R. Sundrum).



- Geometry of extra dimension generates hierarchy exponentially!

$$\Lambda_{\text{TeV}} \sim M_{\text{Planck}} e^{-kL}$$

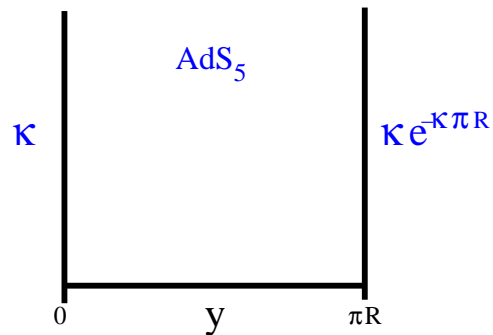
with k the curvature

Warped Extra Dimensions

- Warped 5D metric in RS

$$ds^2 = e^{-2\kappa|y|} \eta^{\mu\nu} dx_\mu dx_\nu + dy^2$$

- Compactified on S_1/Z_2 with $L = \pi R$



and $k \lesssim M_P$, AdS_5 curvature.

- For $kR \simeq (11 - 12)$

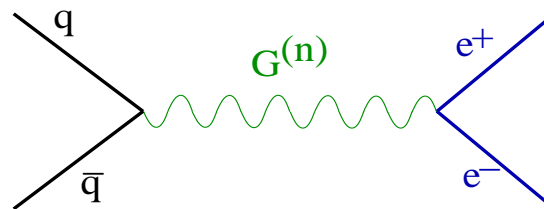
$$\longrightarrow \kappa e^{-\kappa\pi R} \simeq O(\text{TeV}).$$

Warped Extra Dimensions

If only gravity propagates in the bulk (SM fields on TeV brane):

⇒ Kaluza-Klein graviton tower

- Zero-mode graviton $G^{(0)}$ localized toward the Planck brane. This is why gravity is weak! $G^{(0)}$ couples to SM fields as $1/M_P^2$
- First few KK graviton excitations localized toward TeV brane
→ They couple strongly (as $(1/\text{TeV})^2$) to fields there.
E.g.: Drell-Yan at hadron colliders



- In original proposal, *only gravity* propagates in 5D bulk.
- **RS** is a solution of the hierarchy problem. But origin of **EWSB**?
And flavor ? ...
- Allowing gauge fields and matter to propagate in the bulk opens many possibilities: models of EWSB, GUTs, flavor, ...
- The 5D mass of a bulk fermion \Rightarrow *localization* of zero-mode.
- If Higgs remains on TeV brane:
Fermions localized toward TeV brane are more massive
Fermions localized toward the Planck brane are lighter

\Rightarrow **Fermion Geography**

- Gauge Fields in the 5D bulk: KK decomposition in 4D (for $A_y = 0$ gauge):

$$A_\mu(x, y) = \frac{1}{\sqrt{2\pi R}} \sum_{n=0}^{\infty} A_\mu^{(n)}(x) \chi^{(n)}(y) ,$$

Zero-mode $A_\mu(x)^{(0)}$ + KK tower of massive gauge bosons for $n > 0$, with masses

$$m_n \simeq (n - O(1)) \times \pi \kappa e^{-\kappa \pi R}$$

I.e. for appropriate choice of κR 1st KK excitations are $O(\text{TeV})$.

- 1st KK excitations have $\chi^{(n)}(y)$ localized toward TeV fixed point.
- The Gauge Symmetry usually either is or embeds the SM:
e.g. $SU(3)_c \times SU(2)_L \times SU(2)_R \times U(1)_X$ (restores custodial symm.)

- Fermion Fields in the bulk: 5D fermion field KK decomposition

$$\Psi_{L,R}(x, y) = \frac{1}{\sqrt{2\pi R}} \sum_{n=0} \psi_n^{L,R}(x) e^{2\kappa|y|} f_n^{L,R}(y)$$

- 5D fermion bulk mass term \longrightarrow localization of fermion fields:

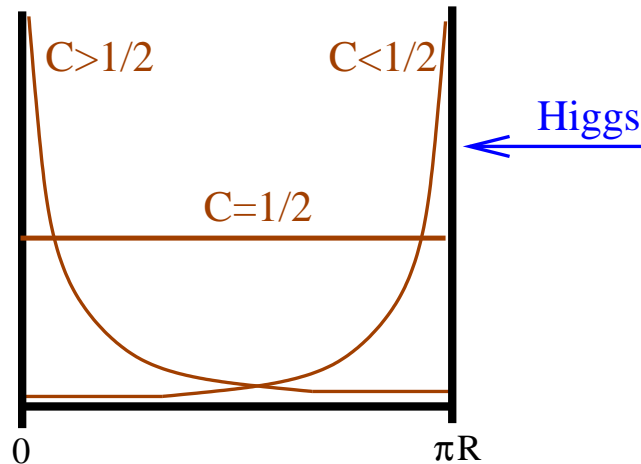
$$S_f = \int d^4x dy \sqrt{-g} \left\{ \dots - c \kappa \bar{\Psi}(x, y) \Psi(x, y) \right\} ,$$

with $c \simeq O(1)$.

- \Rightarrow Fermion zero-modes can be localized by choosing c :

$$f_0^{R,L}(y) = \sqrt{\frac{\kappa\pi R (1 \pm 2c)}{e^{\kappa\pi R(1 \pm 2c)} - 1}} e^{\pm c \kappa y}$$

- $O(1)$ flavor breaking in bulk can generate fermion mass hierarchy:

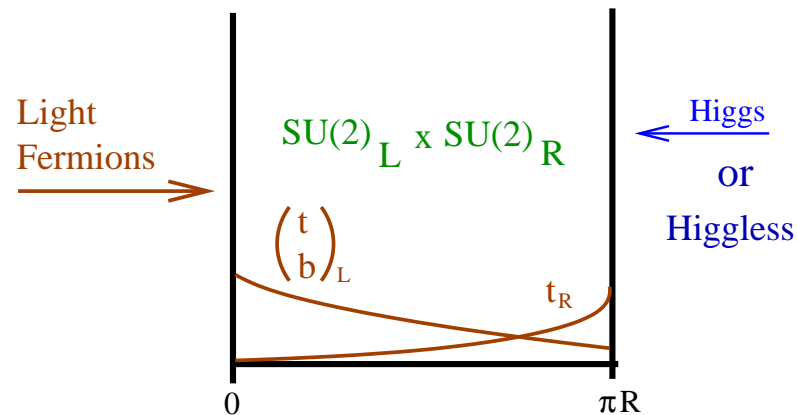


Fermions localized toward the TeV brane can have larger Yukawas,
 Those localized toward the Planck brane have highly suppressed
 ones.

- But fermions at $\simeq \pi R \Rightarrow$ strong couplings to 1st KK gauge bosons!
 E.g: 3rd generation quarks might have large couplings \rightarrow flavor
 violation.

Electroweak Symmetry Breaking and Warped Extra Dimensions

- EWSB by Higgs in the TeV brane. But Electroweak Precision Constraints \Rightarrow bulk gauge symmetry must be $SU(2)_L \times SU(2)_R$ to restore custodial symmetry or T is too large. (Agashe, Delgado, Sundrum '03).



- Can we do without the Higgs ? Higgsless EWSB given by the orbifold boundary conditions (Csaki, Grojean, Murayama, Terning '03). Unitarity is restored by KK gauge bosons. (Chivukula, Dicus, He '01).

Higgsless EWSB and Warped Extra Dimensions

- It is possible to build theories of EWSB and Flavor without a Higgs.
(Csaki, Grojean, Pilo, Terning '03; Nomura '03)
- EW Precision constraints: S tends to be large in AdS/CFT corresp.
→ large N : (Barbieri, Pomarol, Rattazzi '03; Burdman, Nomura '03)

$$S \sim \frac{N}{\pi} = 16\pi \frac{v^2}{m_{KK}^2}$$

- Alternatives:
 - Large m_{KK} , → small N . Leads to strongly coupled “KK” sector. Technicolor-like.
 - Find source of negative S . E.g. kinetic terms on TeV brane (Cacciapaglia, Csaki, Grojean, Terning '04). Allows for low m_{KK} .
But generally constrained by combination of S, T (Chivukula, He, Kurachi, Simmons, Tanabashi '04)

Higgsless Warped Extra Dimensions - Signals

- If cancellation of S is naturally achieved, KK gauge bosons can be rather light and narrow.
- Alternatively, dynamics may be strongly coupled at $O(\text{TeV})$ scale. No spaced resonances, but broad enhancements in cross sections.
- Flavor violation at tree level \rightarrow Potentially rich array of deviations: Effects of KK gluons in CP asymmetries in $B \rightarrow \phi K_s$, $B \rightarrow \pi K_s$, B_s mixing, \dots (Burdman '03);
Z mixing with KK excitations in $b \rightarrow sl^+l^-$ (Burdman, Nomura '03; Agashe, Perez, Soni '04); etc.
- Conclusion: Getting back to strong dynamics at the TeV scale. Can the extra dimensional picture help ?

Conclusions/Outlook

- The Standard Model is an effective description of **EWSB** below the **TeV** scale.
- Physics beyond the SM must be introduced at the **TeV** scale: the SM with a sole Higgs is not radiatively stable.
- The **LHC** will thoroughly explore this energy scale. Starting in 2007!
- What is the solution to the hierarchy problem ?
Supersymmetry, Strong dynamics, Extra dimensions, \dots X ?
- Connections with: Fermion Masses, Grand Unified Theories, Dark Matter. Maybe even Dark Energy.
- Complementarity with Astrophysics/Cosmology/Low energy precision tests.